

# Two stochastic approximations of the chemical source term in the PDF transport equation



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# **Probability Density Function**

## What describes the PDF ?

Consider the velocity  $\underline{U}$ , the scalars  $\underline{\psi}$  (chemical species and enthalpy), turbulent frequency  $\omega$ , and the position  $\underline{x}$  to be a random vector. The quantity

 $f(\underline{V},\underline{\psi},\theta,\underline{x},t)dVd\psi d\theta dx$ 

describes the probability to find at time t the random vector in the interval

 $\underbrace{V} \in [\underline{V}, \underline{V} + d\underline{V}]$   $\underbrace{\psi} \in [\underline{\psi}, \underline{\psi} + d\underline{\psi}]$   $\underbrace{x} \in [\underline{x}, \underline{x} + d\underline{x}]$   $\theta \in [\theta, \theta + d\theta]$ 

# Submodels used

- velocity transport simplified Langevin equation
- scalar transport IEM interaction by exchange with the mean
- turbulent frequency transport Van Sloten, Jayesh, Pope



# **PDF Transport Equation**

### Mean values are obtained

$$\begin{split} \tilde{Q}(\underline{x}) &= \frac{\langle \varrho Q \rangle}{\langle \varrho \rangle} \\ &= \frac{1}{\langle \varrho \rangle} \int_{\underline{V}} \int_{\underline{\psi}} \int_{\theta} Q(\underline{V}, \underline{\psi}, \theta, \underline{x}) \varrho(\underline{\psi}) f(\underline{V}, \underline{\psi}, \theta, \underline{x}) d\underline{V} d\underline{\psi} d\theta \end{split}$$

## Joint velocity-frequency-composition PDF Transport Equation





# **Numerical Solution**

## **Mass Density Function**

 $\mathcal{F}(\underline{U},\underline{\psi},\theta,\underline{x},t)=\varrho(\underline{\psi})f(\underline{U},\underline{\psi},\theta,\underline{x},t)=\langle\varrho\rangle\tilde{f}(\underline{U},\underline{\psi},\theta,\underline{x},t)$ 

### Particle System

 $(\Delta m, (\underline{U}^{(n)}(t), \underline{\varphi}^{(n)}(t), \omega^{(n)}(t), \underline{X}^{(n)}(t))^T) \qquad n = 1, ..., N$ 

#### **Empirical PDF**

$$\mathcal{F}_N(\underline{U},\underline{\psi},\theta,\underline{x}) = \Delta m \sum_{n=1}^N \delta(\underline{U}^{(n)} - \underline{V}) \delta(\underline{\varphi}^{(n)} - \underline{\psi}) \delta(\omega^{(n)} - \theta) \delta(\underline{X}^{(n)} - \underline{x})$$

Weak Convergence !

$$\langle \mathcal{F}_N \rangle / \Delta m \sum_{n=1}^N \langle \delta(\underline{X}^n - x) \rangle \xrightarrow{N \to \infty} \mathcal{F} / \langle \varrho \rangle = \tilde{f}$$

#### **Dynamics of the Particle System**

$$\begin{split} d\underline{U}^{(n)} &= c^{U}(\underline{U}^{(n)}, \partial_{x} \langle p \rangle, \langle \varrho \rangle, \widetilde{\omega}, \widetilde{u}) dt + \sqrt{d^{U}(\underline{U}^{(n)}, \widetilde{\omega}, \widetilde{u}) d\underline{W}} \\ d\underline{\varphi}^{(n)} &= M_{k}(\underline{\varphi}^{(n)}, \widetilde{\varphi}) dt + \underline{R}(\underline{\varphi}^{(n)}) dt \\ d\theta^{(n)} &= c^{\omega}(\omega^{n}, \widetilde{\omega}, \widetilde{u_{i}u_{i}}) dt + \sqrt{d^{\omega}(\theta, \widetilde{\omega}, \widetilde{u_{i}u_{i}})} dW \\ d\underline{X}^{(n)} &= -\underline{U}^{(n)} dt \end{split}$$



# **Numerical Solution**

## Algorithm

- Initialise Particles
  - uniform distribution in space
  - joint normal distribution for velocities
  - gamma distribution for frequency
- Perform Time Step
  - choose time step according to a CFL condition
  - solve SDE system according to the discretised time step
- Estimate Means
  - particle mesh method (cloud in cell)
  - nonparametric kernel estimation
- Correction Algorithm
  - correct particle positions (continuity)
  - correct velocities (zero divergence, mass flux) (mean pressure)
  - calculate mean density from particles
- Next Time Step



# Challenges

## **Computational Time**

- Administration of Particles advanced data structures and sorting algorithms
- Evaluation of Chemistry
  - QSSA
  - ILDM
  - ISAT (ISAT enhanced)
  - PRISM
  - ANN
- Advantages
  - Faster evaluation of the chemical source term
- Disadvantages
  - Usually a tabulation procedure is required
  - High dimension of the parameter space

## Trade off time and storage requirements, and errors



# Weak Chemistry

Particles converge only in weak sense to the PDF

Dynamics of the scalar component of a particle



We need only "stochastic trajectories"! Can we construct a "stochastic trajectory" ? How efficient is this approach?



# The PaSPFR Model

The Cauchy Problem

 $\frac{\partial}{\partial t}f(t,\psi) - \sum_{k=1}^{S+1} \frac{\partial}{\partial \psi_k} \left( \left[ R_k(\psi) + C M_k(f(t,\psi)) \right] f(t,\psi) \right) = 0,$ 

 $f(0,\psi) = f_0(\psi) \,.$ 

The components of the vector  $\psi$  correspond to the mass fractions  $Y_1, \ldots, Y_S$  of several chemical species and to temperature T.

#### The Mixing Model

We use the IEM mixing model.

$$M_k(g,\psi) = \psi_k - \int \dots \int x_k g(x) \, dx_1 \dots dx_{S+1}, \qquad k = 1, \dots, S+1$$

#### The Chemistry Model

The terms  $R_k$  are given by the **reaction mechanism** of the system, which consists of a number I of elementary chemical reactions,

$$(\nu_{\alpha,1},\ldots,\nu_{\alpha,S}) \longleftrightarrow (\nu_{\alpha,1}^*,\ldots,\nu_{\alpha,S}^*), \qquad \alpha=1,\ldots,I$$

The stoichiometric coefficients  $\nu_{\alpha,i}$  and  $\nu_{\alpha,i}^*$  of the species *i* in reaction  $\alpha$  are non-negative integer values.



# **Particle Method**

System of Weighted Particles

$$(w^{(i)}, \Psi^{(i)}(t))$$
  $i = 1, ..., N$ 

where

$$\sum_{i=1}^N w^{(i)} = 1$$

#### **Stochastic Trajectories**

$$\frac{d}{dt}\Psi_k^{(i)}(t) = -R_k(\Psi^{(i)}(t)) - C\left[\Psi_k^{(i)}(t) - \sum_{j=1}^N w^{(j)}\Psi_k^{(j)}(t)\right]$$

where  $k = 1, \ldots, S + 1$ , and  $\Psi^{(i)}(0)$ ,  $i = 1, \ldots, N$ , are independent and distributed according to  $f_0$ .

#### Splitting approach

The **high-dimension** requires a splitting approach on a time interval  $[s, s + \Delta t]$  to decouple the effects of reaction and mixing.



# **Particle Method**

## **Reaction step**

 $\frac{d}{dt}\,\tilde{\Psi}_{k}^{(i)}(t) = -R_{k}(\tilde{\Psi}^{(i)}(t))\,, \qquad t \ge s\,, \qquad \tilde{\Psi}_{k}^{(i)}(s) = \Psi_{k}^{(i)}(s)$ 

#### Mixing step

$$\Psi^{(i)}(s + \Delta t) = e^{-C\Delta t} \tilde{\Psi}^{(i)}(s + \Delta t) + \left[1 - e^{-C\Delta t}\right] \sum_{j=1}^{N} w^{(j)} \tilde{\Psi}^{(j)}(s + \Delta t)$$

#### Algorithm

- 0. Determine the state of the system of particles at time 0 according to the initial density  $f_0\,.$
- 1. Given the state of the system at time s, each particle is moved according to the **reaction step**.
- 2. Given the system  $\tilde{\Psi}^{(1)}(s + \Delta t), \ldots, \tilde{\Psi}^{(N)}(s + \Delta t)$ , each particle is moved according to **mixing step**.
- 3. If time exceeds termination time then STOP. Otherwise go to Step 1.



# The Chemistry Step 1

# The Algorithm

## Step 0

Fix n > 0, set t = 0 and initialise the state vector  $\Psi$ .

## Step 1

Choose a component index  $\boldsymbol{k}$  according to the probability

$$P_k = rac{R_k}{\pi}, \qquad ext{where} \qquad \pi = \sum_{j=1}^{S+1} R_j$$

is called waiting time parameter.

# Step 2

Perform a jump according to

$$\Psi \mapsto J_k(\Psi) = \left(\Psi_1, \dots, \Psi_k - \frac{S}{n} \operatorname{sign} R_k, \dots, \Psi_{S+1}\right),$$

where k is the index chosen in the previous step.

# Step 3

Wait an exponentially distributed time  $\tau$  with waiting time parameter  $\pi$ . That means, advance the time  $t \mapsto t + \tau$  such that the waiting time  $\tau$  is distributed according to

$$\operatorname{Prob}(\tau \ge u) = \exp(-u\pi) \qquad \forall u \ge 0.$$

If t exceeds the splitting step  $\Delta t$  then stop, otherwise go to Step 1.



# The Chemistry Step 2

### A stochastic sub-particle system

$$\Psi = \left( N_1^{(n)}(t), \dots, N_S^{(n)}(t), T^{(n)}(t) \right),$$

The number of sub-particles approximation parameter.

$$n = \sum_{j=1}^{S} N_j^{(n)}(0)$$

### Algorithm

#### **Step 0** The **initial state** of the system is determined as

 $N_j^{(n)}(0) = n X_j(s), \qquad j = 1, \dots, S, \qquad T^{(n)}(0) = T(s),$ 

## Step 1

Calculate the waiting time  $\tau$ 

$$\operatorname{Prob}(\tau \geq u) = \exp(-u\,\pi(x))\,, \qquad u \geq 0\,,$$

where

$$\pi(x) = \sum_{\alpha=1}^{I} |Q_{\alpha,f}(x) - Q_{\alpha,r}(x)|,$$

and  $Q_{\alpha,f}$  and  $Q_{\alpha,r}$  are, respectively, the forward and reverse reaction rates. The procedure stops when time  $t+\tau$  exceeds the splitting step  $\Delta t$ .



# The Chemistry Step 2

## Step 2

At the moment  $t + \tau$ , a particular reaction is chosen according to the **reaction probabilities** 

$$P_{\alpha}(x) = \frac{|Q_{\alpha,f}(x) - Q_{\alpha,r}(x)|}{\pi(x)}, \qquad \alpha = 1, \dots, I.$$

#### Step 3

Finally, the process jumps into the state

$$J_{\alpha}(x) = \begin{cases} J_{\alpha,f}(x) , & \text{if } Q_{\alpha,f}(x) \ge Q_{\alpha,r}(x) ,\\ J_{\alpha,r}(x) , & \text{otherwise} , \end{cases}$$

where

$$J_{\alpha,f}(x) = (x_1 - \nu_{\alpha,1} + \nu_{\alpha,1}^*, \dots, x_S - \nu_{\alpha,S} + \nu_{\alpha,S}^*, x_{S+1} + \Delta T_{\alpha,f}(x))$$

and

$$J_{\alpha,r}(x) = (x_1 - \nu_{\alpha,1}^* + \nu_{\alpha,1}, \dots, x_S - \nu_{\alpha,S}^* + \nu_{\alpha,S}, x_{S+1} + \Delta T_{\alpha,r}(x)).$$

The temperature step is defined as

$$\Delta T_{\alpha,f}(x) = -\frac{\sum_{k=1}^{S} H_k(x_{S+1}) \left[\nu_{\alpha,k}^* - \nu_{\alpha,k}\right]}{\sum_{k=1}^{S} C_k(x_{S+1}) x_k},$$

$$\Delta T_{\alpha,r}(x) = -\frac{\sum_{k=1}^{S} H_k(x_{S+1}) \left[\nu_{\alpha,k} - \nu_{\alpha,k}^*\right]}{\sum_{k=1}^{S} C_k(x_{S+1}) x_k}$$

Then the procedure returns to Step 1.



# **Approximations**

More than 90% of the CPU time is required for updating  $P_{lpha}$ 

### **Combine two strategies:**

- Update  $C_p, H, V, k_{lpha}$  only if T exceeds a given  $\Delta T$
- Update reactions rates not at every reaction event For a chosen constant *ac*,

 $n_{ac}$  is the number of reactions with  $\pi P_{\alpha} \geq \frac{ac}{t_{stop}}$ .

Table 1: The number of sub-particles and  $\Delta T$ .

number of sub-particles	$\Delta T(K)$	ас
$1.0 \times 10^{3}$	10	-
$5.0 \times 10^{3}$	10	-
$1.0 \times 10^{4}$	10	-
$5.0 \times 10^4$	1	2000
$1.0 \times 10^{5}$	1	2000

Table 2:  $n_{ac}$  and the frequency to update reaction rates.

n <sub>ac</sub>	The frequency to	
	update reaction rates	
greater than 50	50	
greater than 25	25	
greater than 10	10	
others	1	

## **Problem description**

non-premixed combustion of n-heptane in hot air

( Detailed mechanism (Westbrook group) 107 chemical species and 808 reversible reactions)

	Fuel	Air	
	$\dot{n}_1$	$\dot{n}_2$	
	2 mol/min	52 mol/min	
X(nC7H16)	1	0	
X(N2)	0	0.7900	
X(O2)	0	0.2100	

Table 3: Initial concentrations and flow rates.

## initial condition approximation

 $\begin{array}{lll} \mbox{Fuel: } 50 \mbox{ particles of weight } w^{(i)} = 2.36 \times 10^{-3}, & i = 1,...,50\,, \\ \mbox{Air: } 50 \mbox{ particles of weight } w^{(i)} = 1.764 \times 10^{-2}, & i = 51,...,100\,, \end{array}$ 





## Splitting error

- Difference in mean temperature at the end of the simulation.
- The CPU time grows linearly with the number of splitting steps.
- From the above study we chose the time step  $\Delta t = 4.0 \times 10^{-4}$ .

Table 4: Splitting error of the temperature at time t = 0.01s.

$\Delta t$	error	
$2.5 \times 10^{-3}$	1249.5	
$1.25 \times 10^{-3}$	417.40	
$6.25 \times 10^{-4}$	26.500	



**Figure 1:** Time evolution of the empirical mean of (a) temperature and (b) OH mass fraction with confidence bounds for 30 repetitions.





**Figure 2:** The deterministic chemistry is compared with the stochastic chemistry for three different approximation parameters. The time evolution of the empirical mean of the following quantities is displayed: (a) temperature. (b) OH mass fraction. (c) n-heptane mass fraction. (d) methane mass fraction.



## Comparison of stochastic and deterministic chemistry

- The **empirical mean** of temperature, OH mass fraction, nheptane mass fraction, and methane mass fraction is used for discussion.
- On the basis of the magnitude of fluctuations one can consider the result for n = 1.0 × 10<sup>4</sup> to be sufficiently accurate. For n = 5.0 × 10<sup>3</sup> the most important characteristics of the ignition are captured.
- **CPU-time** For  $n = 10^4$  the algorithm with stochastic chemistry is about 100 times faster than the algorithm with deterministic chemistry, while it provides a comparable accuracy. In case of  $n = 5. \times 10^3$  sub-particles, we reache a gain factor of more than 200.

method	CPU time [s]	CPU time of single run [s]	speed factor
DASSL	434,759	14,492	1
$5.0 \times 10^4$	22,519	751	19.3
$1.0 \times 10^{4}$	4,260	142	102
$5.0 \times 10^{3}$	2,060	69	211
$1.0 \times 10^{3}$	281	9.4	1,547

**Table 5:** Comparison of computational time of deterministic chemistry andstochastic chemistry approximation.



# Summary and Conclusions

## Stochastic particle method

- New stochastic chemistry approximation
- Partial equilibrium modification
- Incorporation of stochastic chemistry approximation into the PDF transport equation
- Effect of double averaging

## **Numerical Experiments**

- Detailed n-heptane mechanism (106 chemical species and 808 reactions)
- Nonpremixed combustion in a PaSPFR
- Study the splitting error
- Comparison of the stochastic chemistry approach with DASSL
- Stochastic Chemistry Approximation outperformed DASSL for the cases studied.
- More work needs to be done to achieve more significant speed advantage