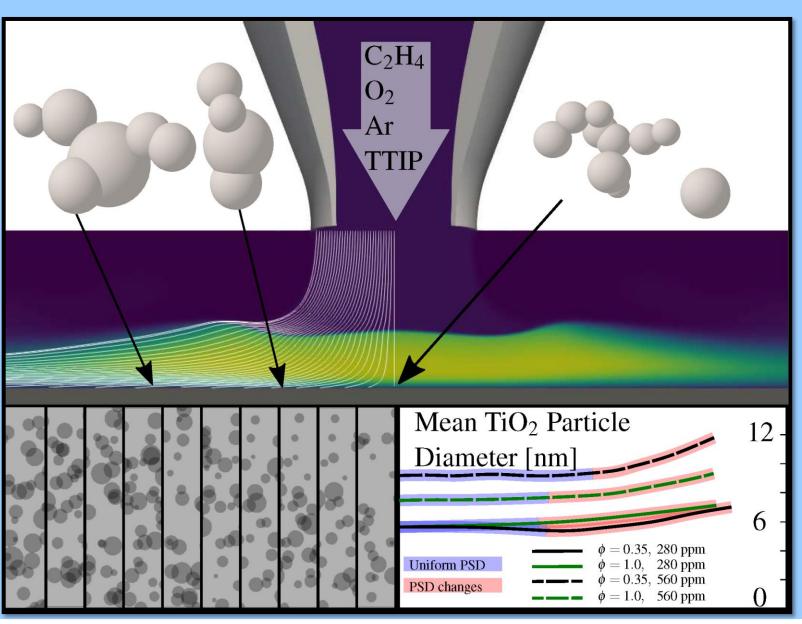
Radial dependence of TiO_2 nanoparticles synthesised in jet-wall stagnation flames

Eric J Bringley, Manoel Y Manuputty, Casper Lindberg, Gustavo Leon, Jethro Akroyd, Markus Kraft

June 25th, 2021





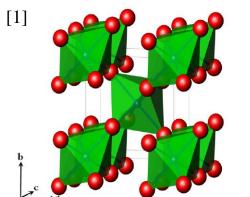
Cambridge Particle Meeting, Cambridge, UK

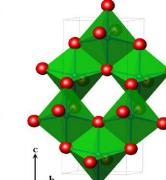


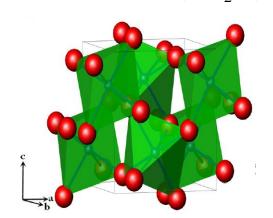
Titania (TiO₂)



Stable Phases (Anatase, Rutile)





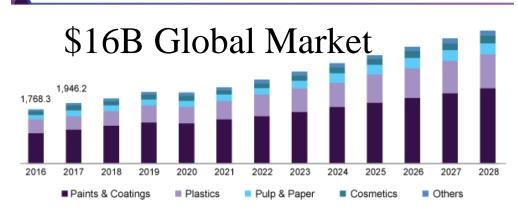


Metastable Phases (TiO₂-II)

Dulux



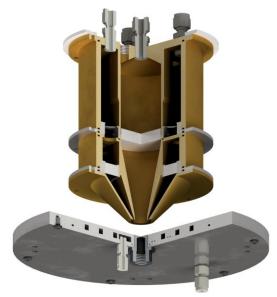
[1] Aravindan et. al (2015) <u>doi: 10.1016/j.mattod.2015.02.015</u>
[2] Wu et. al (2020) <u>doi: 10.1002/smtd.202000928</u>



U.S. TiO2 market size, by application, 2016 - 2028 (USD Million)

Source: www.grandviewresearch.com







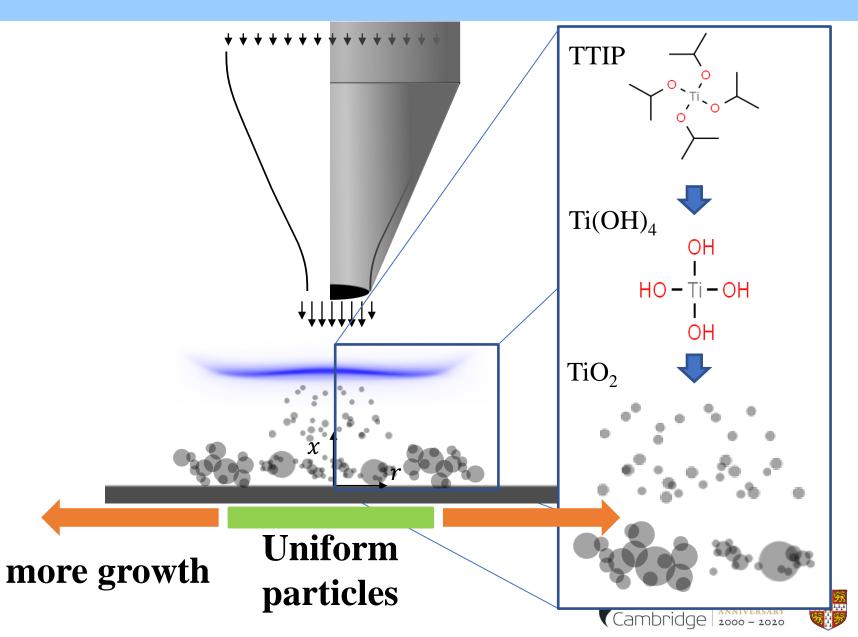
Jet-wall stagnation flames for TiO₂ synthesis

Short residence times produce small particles

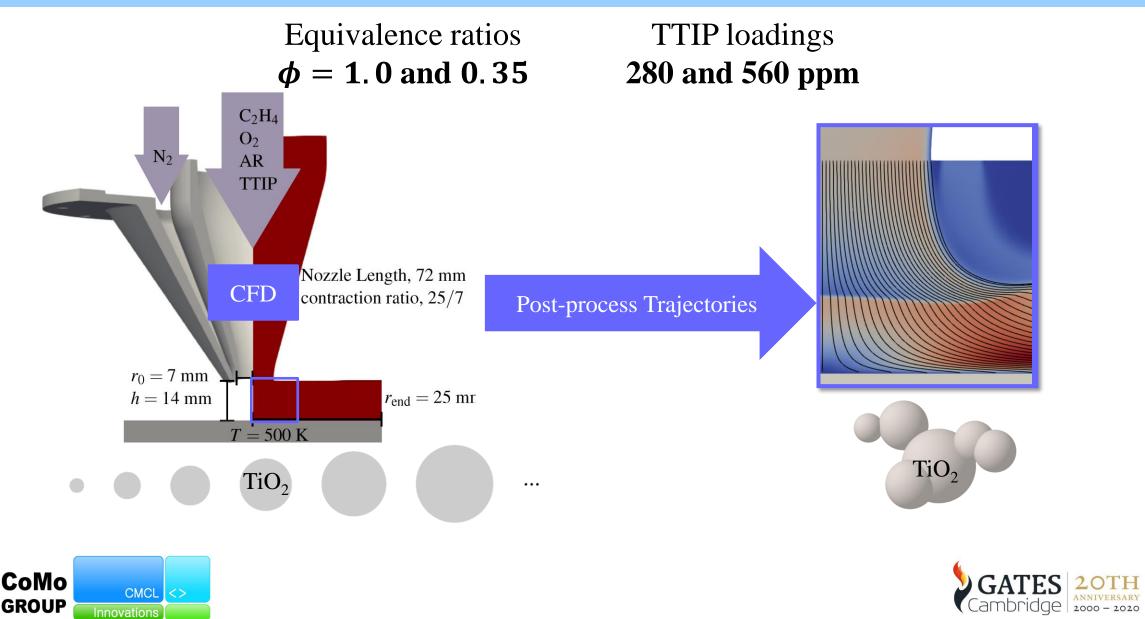
Do the deposited particles change as a function of the radius? If so, where?

 $\approx 1 - 1.5$ nozzle radius



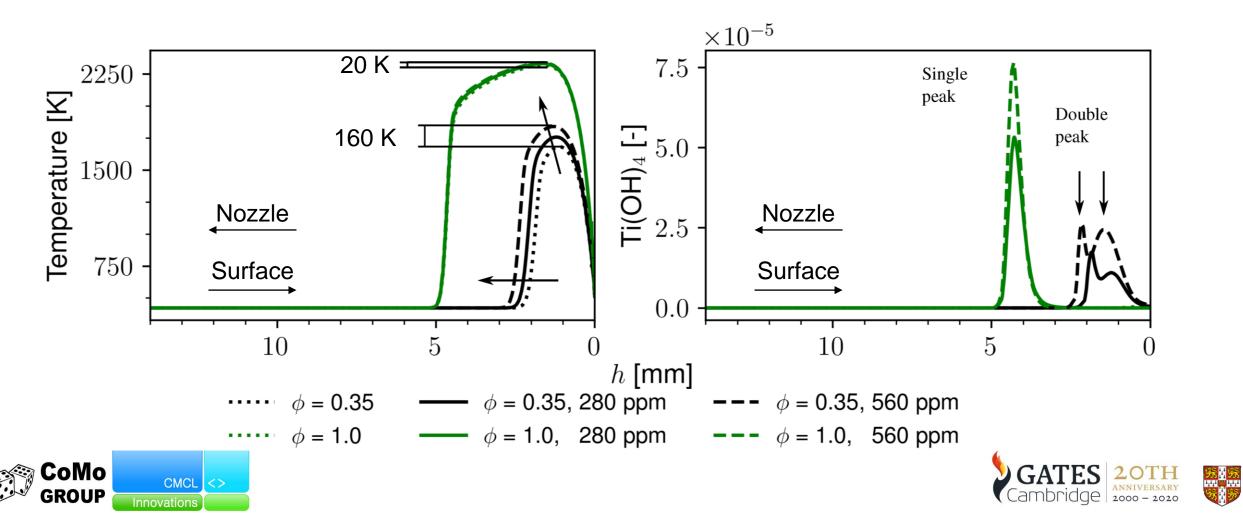


Simulations



Flame centreline data

Two equivalence ratios produce different Ti(OH)₄ profiles



Lagrangian trajectories

As trajectories travel further, 750 the residence time increases. 2250

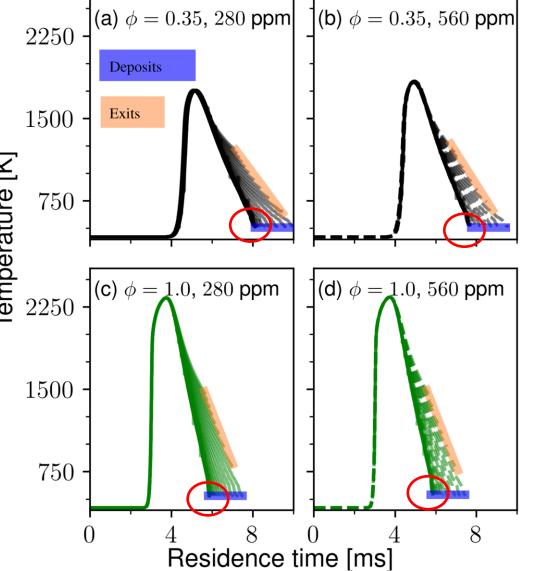
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CoMo

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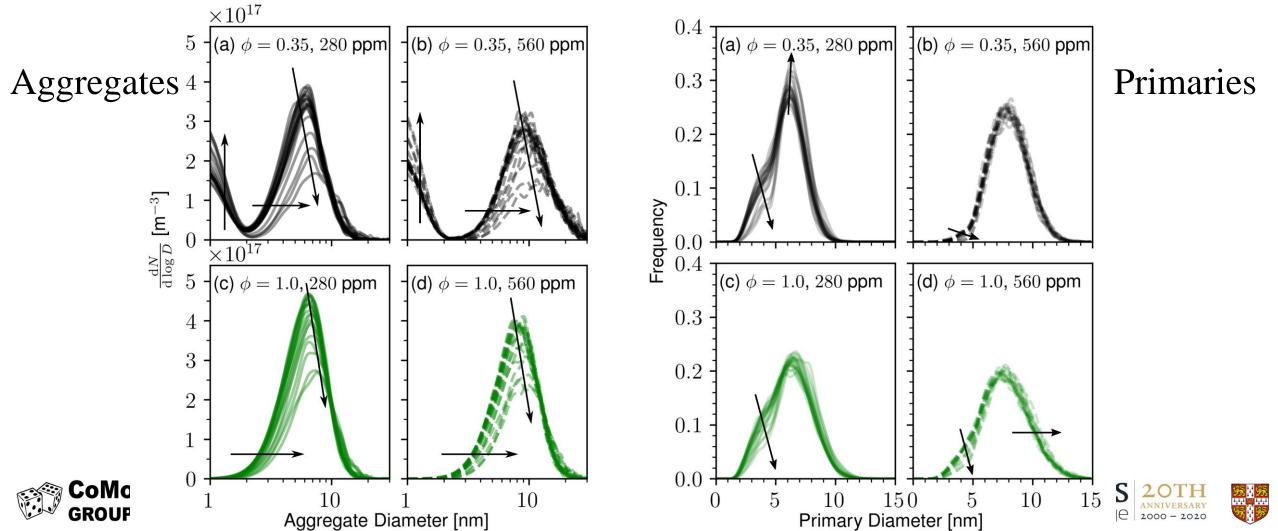
Post-processed with Detailed Particle Model



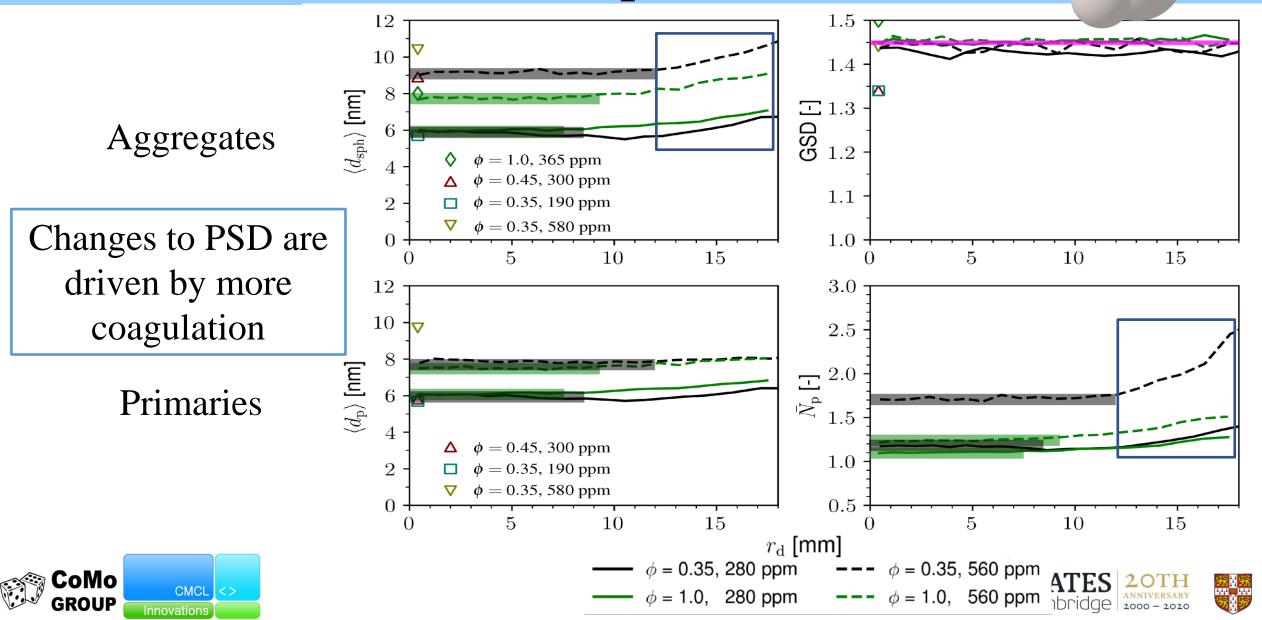
Aggregate and primary size distributions

As trajectories move radially outwards,

the aggregates grow bigger, but primaries remain similar in size



Function of deposition radius



Contributions

Four flames simulated in 2D

 C_2H_4

 O_2

AR

TTIP

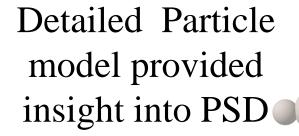
- 500 K

Nozzle Length, 72 mm

contraction ratio, 25/7

 $r_{\rm end} = 25 \, {\rm mr}$

...



PSD differed at ≈ 1.5 nozzle radius due to coagulation 10 $\times 10^{17}$

10

2.5

- 2.0 ~ 1.5

1.0

0.5

$\langle d_{ m sph} angle$ [nm] 5 -(a) $\phi = 0.35, 280$ ppm = 1.0, 365 ppm = 0.45, 300 ppm4 0.35, 190 ppm $\phi = 0.35, 580 \text{ ppm}$ 3 5 3.0 $\mathbf{2}$

15

15

10

10



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Check performance of synthesized particles is radially uniform



 $r_0 = 7 \text{ mm}$

h = 14 mm



Team



Eric Bringley



Gustavo Leon

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Manoel Manuputty



Jethro Akroyd



Casper Lindberg



Markus Kraft





$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho U) = 0;$$

$$\frac{\partial \rho U}{\partial t} + \nabla \cdot (\rho U U) = -\nabla p - \mu \left(\nabla U + (\nabla U)^{\mathsf{T}} \right);$$

$$\frac{\partial \rho Y_i}{\partial t} + \nabla \cdot (\rho U Y_i) + \nabla \cdot (\rho V_i^c Y_i) = \dot{\omega}_i;$$

$$c_p \frac{\partial \rho T}{\partial t} + c_p \nabla \cdot (\rho UT) = \nabla \cdot (\lambda \nabla T) - \left(\rho \sum_{i=1}^N c_{p_i} Y_i V_i^c \right) \cdot \nabla T + \dot{\omega}_{\mathrm{T}}$$

 $\bigotimes_{\text{GROUP}} \operatorname{CoMo}_{\text{GROUP}} = \nabla \cdot \left(\rho U \hat{M}_{j} \right) + \nabla \cdot \left(\rho U \hat{M}_{j} \right) + \nabla \cdot \left(\rho V_{\text{T}} \hat{M}_{j} \right) = \nabla \cdot \left(\rho D_{p_{1}} \nabla \hat{M}_{j-\frac{2}{3}} \right) + \dot{\omega}_{j} \underset{\text{Ge}}{\underset{\text{COV-2020}}{\text{SS}}}$

Closure Models

 $\begin{aligned} & \text{Viscosity} \\ \mu &= \sum_{i=1}^{N} \frac{X_i \mu_i}{\sum_{j=1}^{N} X_j \phi_{ij}} \\ \phi_{ij} &= \frac{1}{\sqrt{8}} \left(1 + \frac{W_i}{W_j} \right)^{-\frac{1}{2}} \left(1 + \left(\frac{\mu_i}{\mu_j} \right)^{\frac{1}{2}} \left(\frac{W_j}{W_i} \right)^{\frac{1}{4}} \right)^2 \end{aligned}$

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Thermal Conductivity

$$\lambda = \frac{1}{2} \left(\sum_{i=1}^{N} X_i \lambda_i + \left[\sum_{i=1}^{N} X_i / \lambda_i \right]^{-1} \right)$$

Diffusion

Binary
$$D_{i,j} = \frac{3}{16} \left(\frac{2N_A k_B^3 T^3}{\pi W_{ij}} \right)^{1/2} \frac{1}{p \sigma_{ij}^2 \Omega^{(1,1*)}}$$
 Mixture-Averaged: $D_i = \frac{1 - Y_i}{\sum_{j \neq i}^N \frac{X_j}{D_{ij}}}$
Diffusion Velocity: $V_i = -D_i \frac{\nabla X_i}{X_i}$, $V_c = -\sum_j^N Y_j V_j$ Corrected Diffusion Velocity: $V_i^c = V_i + V_c$
CoMo



Closure Models

Thermodynamic Properties, JANAF Polynomials $c_{p_i} = \frac{C_{p_i}}{W_i} = \frac{R_g}{W_i} \left(a_{1,i} + a_{2,i}T + a_{3,i}T^2 + a_{4,i}T^3 + a_{5,i}T^4 \right) \qquad c_p = \sum_{i=1}^{N_{sp}} Y_i c_{p_i}$

$$h_{i} = \frac{H_{i}}{W_{i}} = \frac{R_{g}}{W_{i}} \left(a_{1,i}T + \frac{a_{2,i}}{2}T^{2} + \frac{a_{3,i}}{3}T^{3} + \frac{a_{4,i}}{4}T^{4} + \frac{a_{5,i}}{5}T^{5} + a_{6,i} \right) \qquad h = \sum_{i=1}^{N_{sp}} Y_{i}h_{i}$$

Ideal Gas Law

$$pV = nRT$$



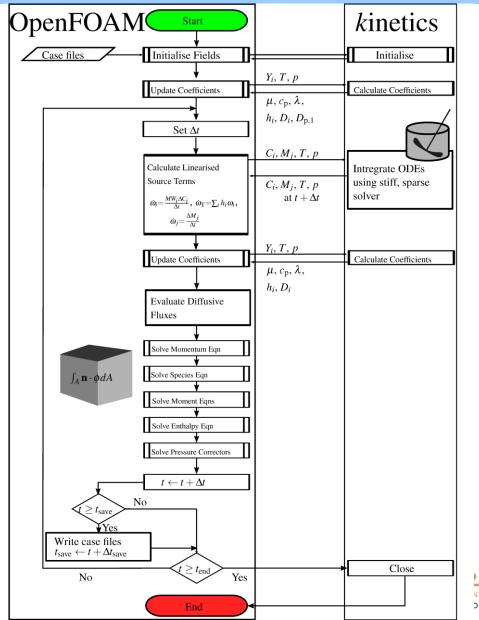


Models and methods: flow

2D Simulations: Navier Stokes Equations, CFD with PISO Alg.

CFD Models: Ideal Gas Law, JANAF, Mixture Avg. Transport, UCSD Chemistry







Models and methods: TiO₂ particles

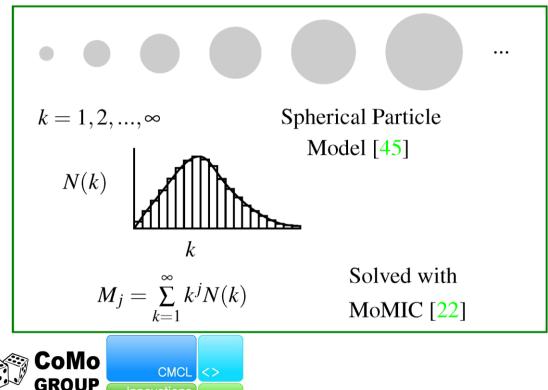
Computational efficiency

vs Physical insight

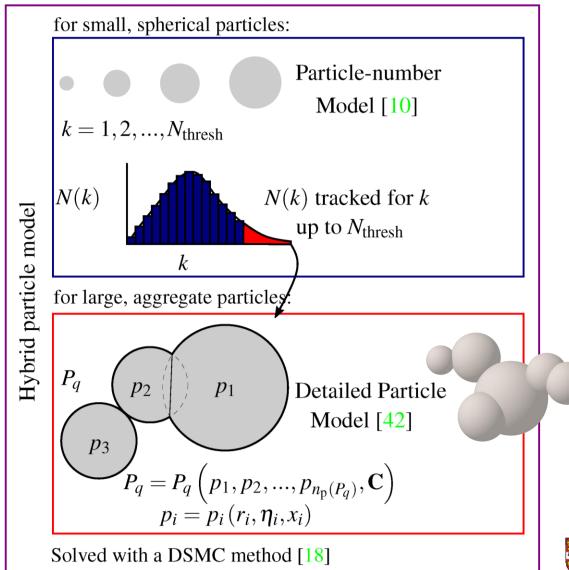
Same gas phase transfer species, Ti(OH)₄ Particle inception, surface growth, and coagulation

Coupled to governing equations

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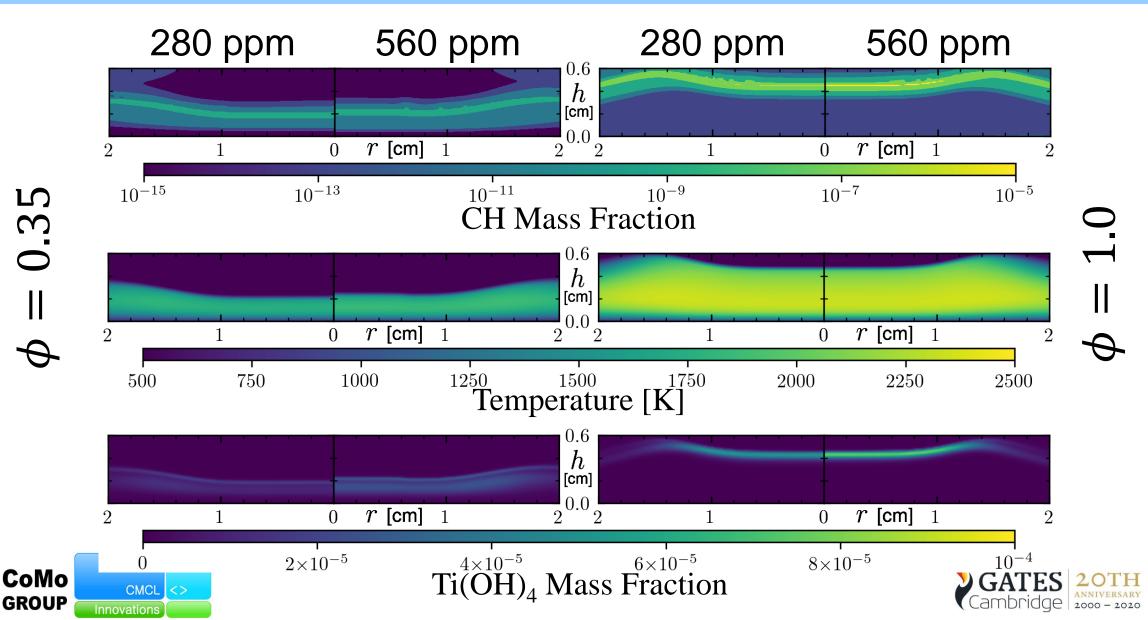


Lagrangian post-processing

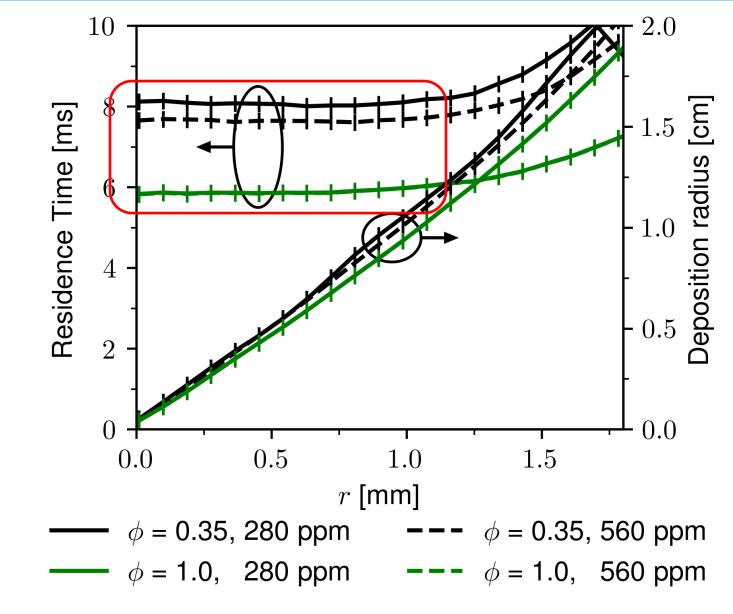




Flame data



Trajectory properties



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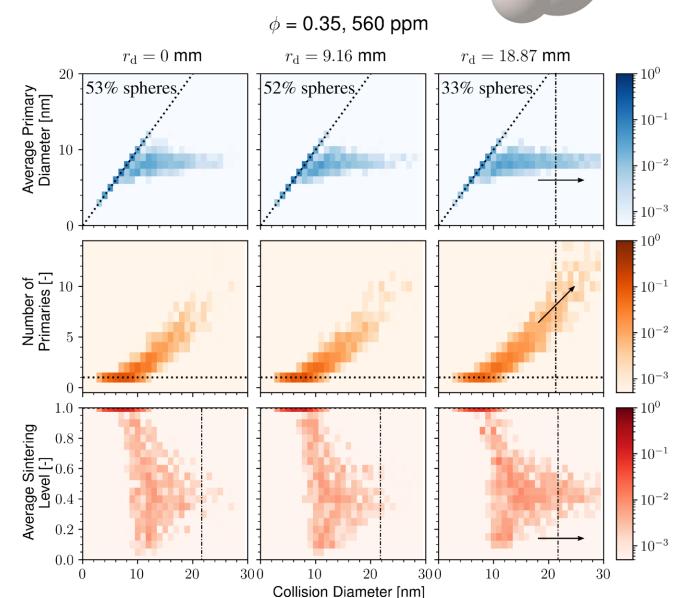


Primaries particle joint property distributions

Average primary diameter does not significantly change

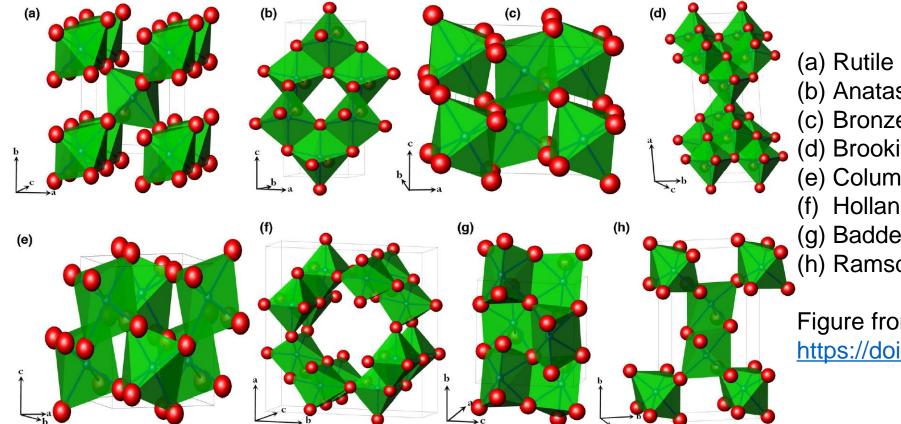
Large particles at large radius have a high number of primaries

Large particles are weakly sintered





Titania crystal phases



(b) Anatase (c) Bronze (d) Brookite (e) Columbite (f) Hollandite (g) Baddeleyite (h) Ramsdellite

Figure from Aravindan et. al (2015), https://doi.org/10.1016/j.mattod.2015.02.015



